

A CONICAL SHELL THEORY OF HYBRID ANISOTROPIC MATERIALS

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ABSTRACT

We start with the unique coordinate system of a conical shell and the generalized Hook's law of a three dimensional anisotropic body which is subjected to 6 different stresses and strains together with 21 elastic coefficients. Applying equilibrium and stress displacement equations and non-dimensionalize all the variables before asymptotically expanding, we are now allowed to perform the asymptotic integration process which will derive the first approximation theory of an anisotropic conical shell. The theory is valid for anisotropic non-homogeneous materials, which can be applied to the layered walls of conical shells with different materials. A choice of small length scale allows a system of differential equations accurately formulated for different loading and boundary conditions. The derived differential equations of homogeneous isotropic material are equivalent to classical theory.

Keywords: Conical Shell, Three Dimensional Anisotropic Element, Non-Homogeneous Anisotropic Materials, Generalized Hook's Law, Boundary Layer, Asymptotic Integration, Classical Shell Theory.

INTRODUCTION

The use of composite materials have been wide spread over last 30 years from the aerospace vehicles to civil engineering structures, the materials possess such characteristics as high strength/density and properly controllable modulus/density ratios and they are generally composed of filaments and matrix materials. The filaments are embedded in the matrix materials to give additional stiffness and high strength.

The cylindrical, spherical and conical shell theories that we develop in our previous articles are common to liquid oxygen storage tanks of outer space rockets, orbital shuttles as well as liquid storage tank and building structures. The authors already developed and formulated almost all of necessary cylindrical and spherical shell theories of multiple mixture (hybrid) anisotropic materials. Details are published in various journals as shown in the References.

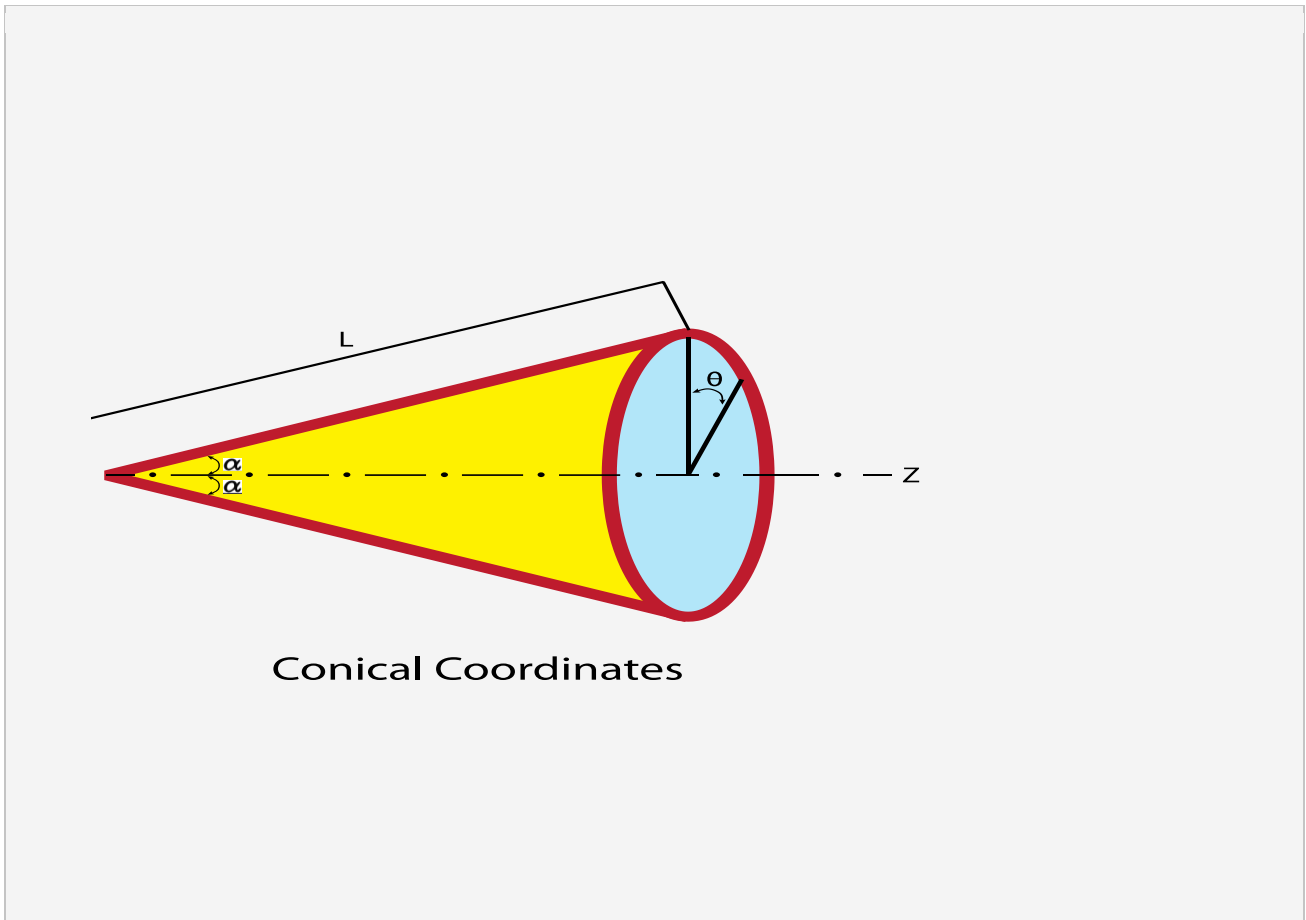


Figure (1) Coordinates

As shown in Figures 2 and 3 the filaments can be arranged arbitrarily to make a composite structure more resistant to loadings. As the mechanical properties of composites vary depending on the direction of the fiber arrangement, it is necessary to analyze them by use of an anisotropic theory. Composite materials are also generally constructed of thin layers, which may have different thickness and different cross-ply angles. The cross-ply angle, γ , is the angle between major elastic axis of the material and reference axis (see Figures 2 and 3). The variation in properties in the direction of the thickness implies non-homogeneity of the material and composite structures must thus be analyzed according to theories, which allow for non-homogeneous anisotropic material behavior.

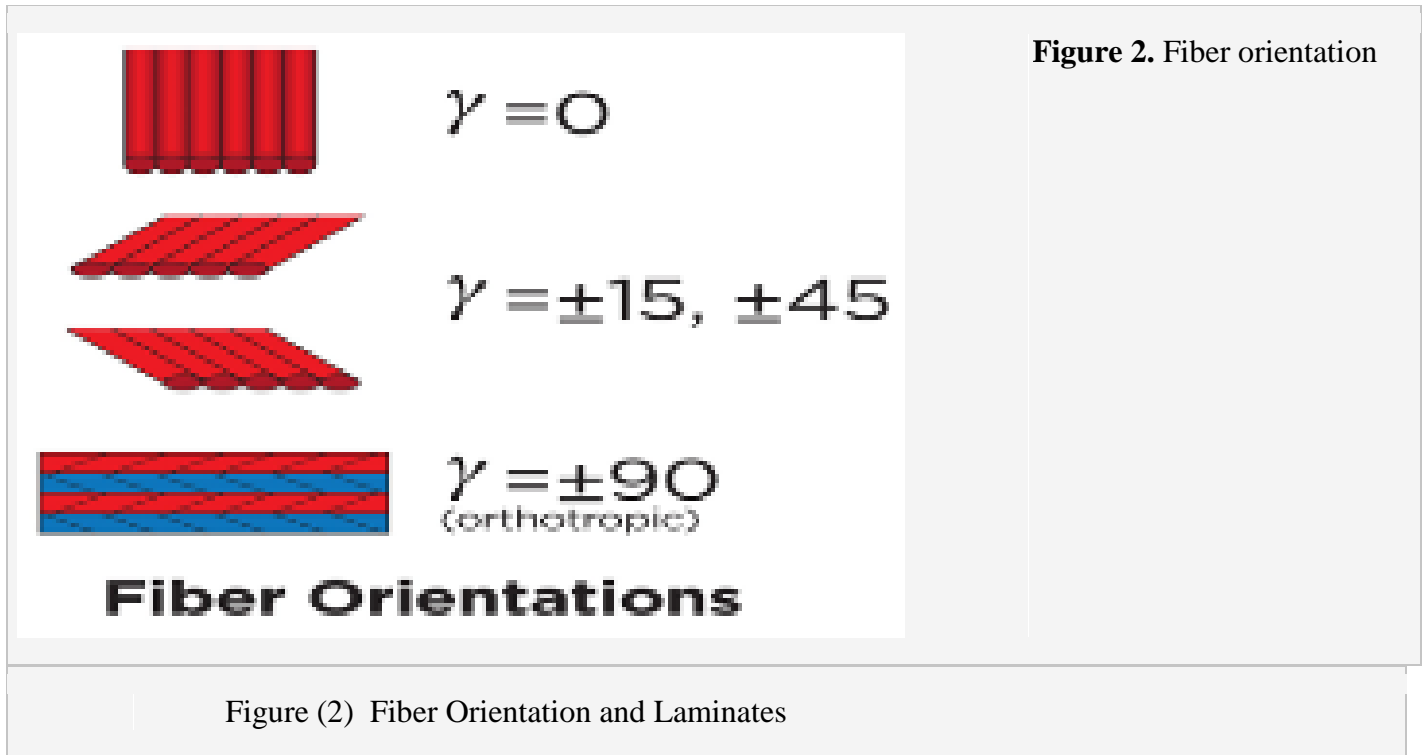


Figure (2) Fiber Orientation and Laminates

According to the exact three-dimensional theory of elasticity, a shell element is considered as a volume element. All possible stresses and strains are assumed to exist and no simplifying assumptions are allowed in the formulation of the theory. We therefore allow for six stress components, six strain components and three displacements. There are thus fifteen unknowns to solve for in a three-dimensional elasticity problem. On the other hand, three equilibrium equations and six strain displacement equations can be obtained for a volume element and six generalized Hook's law equations can be used. A total of fifteen equations can thus be formulated and it is possible to set up a solution for a three-dimensional elasticity problem. It is however very complicated to obtain a unique satisfaction of the above fifteen equations and the associated boundary conditions. This led to the development of various theories for structures of engineering interest. A group of simplifying assumptions that provide a reasonable description of the behavior of thin elastic shells proposed by Love (Reference [1]) has led to the development of the classical shell theories.

Love's theory is based upon the well-known hypotheses of thin shell theory. A detailed description of classical shell theory can be found in various articles. (References [2] to [6]). The system of shell equations derived based on the assumptions is usually regarded as satisfactory for thin homogenous, isotropic shells except for problems associated with certain types of loadings and where the transverse stresses and strains may be important. Being free from the classical assumptions is to apply the asymptotic method to the three-dimensional elasticity equations and thus obtain so-called rational two-dimensional shell theories. Asymptotic methods have at their foundation the desire to obtain a solution that is approximately valid when a physical parameter (or a variable) of the problem which is very small (or very large). The solution is usually of a boundary or initial value problem in powers of a parameter, which either appears explicitly in the original problems or is introduced in some artificial manner. The highlights of using the

asymptotic method to derive shell equations are first to comply with the three-dimensional elasticity theory and incorporates the expansion of stress and displacement components in an asymptotic series. We can then collect the first approximation by taking only the leading term of the expansions.

The derivation of the theories is accomplished by first introducing the shell coordinates and dimensions and yet unspecified characteristic length scales via changes in the independent variables. Next, the dimensionless stresses and displacements are expanded asymptotically by using the thinness of the shell as the expansion parameter. A choice of characteristic length scales is then made, and corresponding to different combination of these length scales, different sequences of systems of differential equations are obtained. Subsequent integration over the thickness and satisfaction of the boundary conditions yields the desired equations governing the formulation of the first approximation theory of non-homogeneous anisotropic cylindrical shells.

General Anisotropic Conical Shell Theory

Based on a non-homogeneous, anisotropic volume element of a conical body with longitudinal, circumferential (angular) and radial coordinates being noted as z , θ and r , respectively, and subjected to all possible stresses and strains (Figure [1]). The cone occupies the space between $a \leq r \leq a + h$ and the edges are located at $L = 0$ and $L = L_0$. Here, a is the inner radius, h the thickness and L the cone length.

Assuming that the deformations are sufficiently small so that linear elasticity theory is valid, the following equilibrium and stress-displacement equations [Equations (2.1) and (2.2)] will govern the problem:

$$\begin{aligned} (r\tau_{rz})_{,r} + \tau_{\theta z,\theta} + \sin \alpha (r \sigma_z)_{,z} &= 0 \\ (r\tau_{r\theta})_{,r} + \sigma_{\theta,\theta} + \sin \alpha (r \tau_{\theta z})_{,z} + \tau_{\theta z} &= 0 \\ (r\sigma_r)_{,r} + \tau_{r\theta,\theta} + \sin \alpha (r \tau_{rz})_{,z} - \sigma_{\theta} &= 0 \end{aligned} \quad (2.1)$$

$$\begin{aligned} U_{z,z} &= S_{11} \sigma_z + S_{12} \sigma_{\theta} + S_{13} \sigma_r + S_{14} \tau_{r\theta} + S_{15} \tau_{rz} + S_{16} \tau_{\theta z} \\ (1/\sin \alpha)(u_{\theta,\theta} + u_r) / r &= S_{21} \sigma_z + S_{22} \sigma_{\theta} + S_{23} \sigma_r + S_{24} \tau_{r\theta} + S_{25} \tau_{rz} + S_{26} \tau_{\theta z} \\ (\sin \alpha) u_{r,z} &= S_{31} \sigma_z + S_{32} \sigma_{\theta} + S_{33} \sigma_r + S_{34} \tau_{r\theta} + S_{35} \tau_{rz} + S_{36} \tau_{\theta z} \\ (1/\sin \alpha) U_{r,\theta} (1/r) + u_{\theta,z} - (1/\sin \alpha) u_{\theta} / r &= S_{41} \sigma_z + S_{42} \sigma_{\theta} + S_{43} \sigma_r + S_{44} \tau_{r\theta} + S_{45} \tau_{rz} + S_{46} \tau_{\theta z} \\ (1/\sin \alpha) U_{z,r} + \sin \alpha u_{r,z} &= S_{51} \sigma_z + S_{52} \sigma_{\theta} + S_{53} \sigma_r + S_{54} \tau_{r\theta} + S_{55} \tau_{rz} + S_{56} \tau_{\theta z} \\ u_{\theta,z} + (1/\sin \alpha) U_{z,\theta} / r &= S_{61} \sigma_z + S_{62} \sigma_{\theta} + S_{63} \sigma_r + S_{64} \tau_{r\theta} + S_{65} \tau_{rz} + S_{66} \tau_{\theta z} \end{aligned} \quad (2.2)$$

In the above, u_r , u_{θ} and u_z are the displacement components in the radial, circumferential and longitudinal directions, respectively but all measured by normal radius of

curvature, therefore must be adjusted by the conical angle α . σ_r , σ_θ , and σ_z are the normal stress components in the same directions and τ_{rz} , $\tau_{\theta z}$ and $\tau_{r\theta}$ are the shear stresses on the θ - z face, r - z face and r - θ face, respectively (Figure 1). A comma indicates partial differentiation with respect to the indicated coordinates. The S_{ij} 's ($i, j = 1, 2, \dots, 6$) in (2.2) are the components of compliance matrix and represent the directional properties of the material. The compliance matrix is symmetric, $S_{ij} = S_{ji}$. Complete anisotropy of the material is allowed for, making 21 independent material constants. We are not allowed to eliminate any of those components since the material properties depend on the manufactures set up, or, in the case of aerospace vehicles, a different gravity environment. The components can be expressed in terms of engineering constants as follows:

$$\begin{aligned}
 S_{ii} &= \frac{1}{E_i}, \quad (i, j = 1, 3) \\
 S_{ij} &= \frac{-\nu_{ij}}{E_i}, \quad (i = 1, 2, j = 2, 3, i \neq j) \\
 S_{44} &= \frac{1}{G_{23}} \\
 S_{55} &= \frac{1}{G_{13}} \\
 S_{66} &= \frac{1}{G_{12}}
 \end{aligned} \tag{2.3}$$

In (2.3), the E_i 's are the Young's moduli in tension along the i - direction and ν_{ij} and G_{ij} are the Poisson's ratio and shear moduli in the i - j face, respectively. Equation (2.2) implies anisotropic property of the material only. Materials to be non-homogeneous, different properties of each layer of the shell, we will allow the material property variation in the radial direction as follows:

$$S_{ij} = S_{ij}(r) \tag{2.4}$$

Unlike the majority of theories, which identify material properties from the beginning, this theory will allow the properties to be, variable, as hybrid theory.

The principal material axes (r', θ', z') in general do not coincide with the body axes of the conical shell (r, θ, z). If the material properties S'_{ij} with respect to material axes specified, then the properties with respect to the body axes are given by the anisotropic transformation equations:

The shell is free from surface traction at its inner surface while the outer surface is subjected to a uniformly distributed tensile force. The boundary condition is then:

$$\begin{aligned} \sigma_r = \tau_{r\theta} = \tau_{rz} &= 0 & (r = a) \\ \sigma_r = p(\theta, z), \tau_{r\theta} = \tau_{rz} &= 0 & (r = a + h) \end{aligned} \quad (2.5)$$

We will find it convenient to work with stress resultants rather than the stresses themselves. These stress resultants, which are forces and moments per unit length, are obtained by integrating with respect to the thickness coordinate. They are:

$$\begin{aligned} N_z &= \int_a^{a+h} \sigma_z (\sin \alpha) \left[1 + \frac{r - a - d}{a + d} \right] dr \\ N_\theta &= \int_a^{a+h} \sigma_\theta (\sin \alpha) dr \\ N_{\theta z} &= \int_a^{a+h} \tau_{\theta z} (\sin \alpha) dr \\ N_{z\theta} &= \int_a^{a+h} \tau_{\theta z} (\sin \alpha) \left[1 + \frac{r - a - d}{a + d} \right] dr \\ M_\theta &= \int_a^{a+h} \sigma_\theta (\sin \alpha) [r - a - d] dr \\ M_z &= \int_a^{a+h} \sigma_z (\sin \alpha)^2 \left[\frac{r - a - d}{a + d} \right] r dr \\ M_{\theta z} &= \int_a^{a+h} \tau_{\theta z} (\sin \alpha) [r - a - d] dr \\ M_{z\theta} &= \int_a^{a+h} \tau_{\theta z} (\sin \alpha)^2 \left[\frac{r - a - d}{a + d} \right] r dr \end{aligned} \quad (2.7)$$

In the above, a is the inner radius of the conical shell and d is the distance from the inner surface to the reference surface where the stress resultants are defined. Note that $N_{\theta z}$ and $N_{z\theta}$ and $M_{z\theta}$ and $M_{\theta z}$ respectively are different. This is due to the fact that the terms of the order of thickness over radius are not neglected compared to one in the integral expressions.

Formulation of a Boundary Layer Theory

As explained in the introduction, a theory of shells is distinguished from the exact three dimensional elasticity formulation by the fact that one of the coordinates is suppressed by the mathematical description. The procedure used here for obtaining the two dimensional thin shell equations is that of the asymptotic integration of the (2.1) and (2.2) describing the cylindrical shell. As a first step for integrating (2.1) and (2.2), we non-dimensionalize the coordinates as follows:

$$X = z/L, \quad y = (r - a)/h, \quad \varphi = \theta / \beta \quad (3.1)$$

$$X = z/L, \quad y = (r - a)/h, \quad 0 = 0 / B$$

where L and L_0 are quantities which are to be determined later.

Next, the compliance matrix, the stresses and deformations are non-dimensionalized by the use of a representative stress level σ , a representative material property S and the shell radius a as follows:

$$\begin{aligned} S_{ij} &= S \bar{S}_{ij} \\ \sigma_z &= \sigma t_{r\theta}, \quad \sigma_\theta = \sigma t_\theta, \quad \sigma_r = \sigma t_r \\ \tau_{r\theta} &= \sigma t_{r\theta}, \quad \tau_{rz} = \sigma t_{rz}, \quad \tau_{\theta z} = \sigma t_{\theta z} \\ u_r &= \sigma a S v_r, \quad u_\theta = \sigma a S v_\theta, \quad u_z = \sigma a S v_z \end{aligned} \quad (3.2)$$

Where the dimensionless displacements and stresses are functions of x, y and φ . These variables, together with their derivatives with respect to x, y and φ , are assumed to be of order unity. The parameters L and ℓ introduced in (3.1) are thus seen to be characteristic length scales for changes of the stresses and displacements in the axial and circumferential directions, respectively.

Consider a small parameter ϵ , $\epsilon < 1$. With respect to an arbitrary domain D of the cylinder, M_1 is said to be of order ϵ^n relative to a second quantity M_2 .

$$M_2 \approx \epsilon^n M_1 \quad (3.3)$$

Where n is an arbitrary between 0 and infinity

If everywhere in D (with the possible exception of some isolated small regions) the relation

$$\epsilon^{n+m} \leq |M_2|/|M_1| \leq \epsilon^{n-m} \quad (3.4)$$

holds for a suitably chosen value of m , $0 < m < 1$. According to this definition, two quantities are of the same order if $n = 0$ in the above, while a quantity is of order unity when $n = 0$ and $M_1 = 1$. Substitution of the dimensionless variables defined by (3.1) and (3.2) into the elasticity equations, (2.1) and (2.2), yields the following dimensionless equations:

$$\begin{aligned} v_{r,y} &= \lambda [\bar{S}_{31}t_z + \bar{S}_{32}t_\theta + \bar{S}_{33}t_r + \bar{S}_{34}t_{r\theta} + \bar{S}_{35}t_{rz} + \bar{S}_{36}t_{\theta z}] / \sin \alpha \\ v_{z,y} + \sin \alpha \lambda \left(\frac{a}{l}\right) v_{r,x} &= \lambda [\bar{S}_{51}t_z + \bar{S}_{52}t_\theta + \bar{S}_{53}t_r + \bar{S}_{54}t_{r\theta} + \bar{S}_{55}t_{rz} + \bar{S}_{56}t_{\theta z}] \\ \lambda \sin \alpha v_{r,0} + \left(\frac{l}{a}\right) (1 + \lambda y) v_{\theta,y} - \left(\frac{l}{a}\right) \lambda v_\theta & \\ = \lambda \left(\frac{l}{a}\right) (1 + \lambda y) [\bar{S}_{41}t_z + \bar{S}_{42}t_\theta + \bar{S}_{43}t_r + \bar{S}_{44}t_{r\theta} + \bar{S}_{45}t_{rz} + \bar{S}_{46}t_{\theta z}] & \quad (3.5) \\ v_{z,x} = \left(\frac{l}{a}\right) [\bar{S}_{11}t_z + \bar{S}_{12}t_\theta + \bar{S}_{13}t_r + \bar{S}_{14}t_{r\theta} + \bar{S}_{15}t_{rz} + \bar{S}_{16}t_{\theta z}] & \\ \left(\frac{a}{l}\right) v_{\theta,0} + (\sin \alpha) v_r = (1 + \lambda y) [\bar{S}_{21}t_z + \bar{S}_{22}t_\theta + \bar{S}_{23}t_r + \bar{S}_{24}t_{r\theta} + \bar{S}_{25}t_{rz} + \bar{S}_{26}t_{\theta z}] & \\ \left(\frac{a}{l}\right) (1 + \lambda y) v_{\theta,x} + \left(\frac{a}{l}\right) v_{z,0} = (1 + \lambda y) [\bar{S}_{61}t_z + \bar{S}_{62}t_\theta + \bar{S}_{63}t_r + \bar{S}_{64}t_{r\theta} + \bar{S}_{65}t_{rz} + \bar{S}_{66}t_{\theta z}] & \end{aligned}$$

$$\begin{aligned} [t_{rz}(1 + \lambda y)]_{,y} + \left(\frac{\lambda a}{l}\right) t_{\theta z,0} + \left(\frac{\lambda a}{l}\right) (1 - \lambda y) t_{z,x} &= 0 \\ [t_{r\theta}(1 + \lambda y)]_{,y} + \left(\frac{\lambda a}{l}\right) t_{\theta,0} + \lambda t_{r\theta} + \left(\frac{\lambda a}{l}\right) (1 + \lambda y) t_{\theta z,x} &= 0 \\ [t_r(1 + \lambda y)]_{,y} + \left(\frac{\lambda a}{l}\right) t_{r\theta,0} + \left(\frac{\lambda a}{l}\right) (1 + \lambda y) t_{rz,x} - \lambda t_\theta &= 0 \end{aligned} \quad (3.6)$$

Where λ is the thin shell parameter, defined as

$$\lambda = h/a \quad (3.7)$$

The parameter λ is representative of the thickness of the conical shell. We will consider only the case of thin shell theory as

$$\lambda \ll 1 \quad (3.8)$$

The dimensionless coefficients S_{ij} of the compliance matrix in general are not all of same order. We therefore assume that they can be expanded in terms of finite sum as:

$$\overline{S}_{ij}(y; \lambda) = \sum_{n=0}^N S_{ij}^{(n)}(y) \lambda^{\frac{n}{2}} \quad (3.9)$$

Where $S_{ij}^{(n)}(y)$ is of order unity or vanish identically. Next, we assume that each displacement components, represented by the generic symbol v and each stress component, represented by the generic symbol t , can be expanded in terms of a power series in $\lambda^{1/2}$

$$\begin{aligned} v(y, x, \phi; \lambda) &= \sum_{m=0}^M v^{(m)}(y, x, \phi) \lambda^{m/2} \\ t(y, x, \phi; \lambda) &= \sum_{m=0}^M t^{(m)}(y, x, \phi) \lambda^{m/2} \end{aligned} \quad (3.10)$$

The $v^{(m)}$ and $t^{(m)}$ are of order unity. No convergence properties are assumed for series (3.10) except only asymptotic validity for λ . That is, if expansions (3.10) are terminated at some power of $\lambda^{1/2}$, the error in using the expansions rather than the exact solutions tends to zero as λ approaches zero.

Length scales L are as yet arbitrary. Their choice, as will be seen in the subjects to follow, determines the type of shell theory to be identified. Based in part on the results obtained for isotropic, homogeneous shells as shown in Reference [6], the developed theory is equivalent to Donnell-Vlasov thin cylindrical shell equation.

The last step of the procedure consists of substituting the expansions for the series and assumed length scales into the dimensionless elasticity equations of stress-displacement and equilibrium described by (3.5) and (3.6). Upon selecting terms of like powers in $\lambda^{1/2}$ on both sides of each equations and requiring that the resulting equations be integrable with respect to the thickness coordinate as well as be capable of identifying the relations for all stresses and displacement

components, we will obtain systems of differential equations. The first system of equations of “thin shell” theory and we will call it the first approximation system. We can, however, obtain stresses and displacements of each layer of thickness coordinate, which can be an advantage of the procedure among others. In the following section, the thin shell theories for the combinations of length scales can be derived.

Shell Theory of Longitudinal and Circumferential Length Scales $(Lh)^{1/2}$

We will now develop a theory of both characteristic length scales same as $(Lh)^{1/2}$, as follows

$$L = (Lh)^{1/2}, \quad 1 = (Lh)^{1/2} \quad (4.1)$$

In the above a is inner radius of conical shell and h is the total thickness of the wall as shown in Figure 1 and Figure 3.

The basis of the above equations are first non-dimensionalize each and all the variables to compare the magnitudes on equal basis and empirical number, which is used in the theory of shells.

The following systems of differential equations are obtained by substituting the characteristic length scales (4.1) into (3.5) and (3.6).

By substituting the asymptotic expansions (3.10) for the displacements and stresses and expansions (3.9) for the compliance matrix, we obtain the following systems of equations representing the first approximation theory.

$$v_{r,y} = \lambda \left[\bar{S}_{31} t_z + \bar{S}_{32} t_\theta + \bar{S}_{33} t_r + \bar{S}_{34} t_{r\theta} + \bar{S}_{35} t_{rz} + \bar{S}_{36} t_{\theta z} \right] / \sin \alpha$$

$$v_{z,y} + \lambda^{\frac{1}{2}} v_{r,x} (\sin \alpha) = \lambda \bar{S}_{51} t_z + \bar{S}_{52} t_\theta + \bar{S}_{53} t_r + \bar{S}_{54} t_{r\theta} + \bar{S}_{55} t_{rz} + \bar{S}_{56} t_{\theta z}$$

$$\lambda (\sin \alpha) v_{r,\phi} + \lambda^{\frac{1}{2}} (1 + \lambda y) v_{\theta,y} - \lambda^{3/2} v_\theta$$

$$= \lambda^{3/2} (1 + \lambda y) \bar{S}_{41} t_z + \bar{S}_{42} t_\theta + \bar{S}_{43} t_r + \bar{S}_{44} t_{r\theta} + \bar{S}_{45} t_{rz} + \bar{S}_{46} t_{\theta z}$$

$$v_{z,x} = \lambda^{\frac{1}{2}} \bar{S}_{11} t_z + \bar{S}_{12} t_\theta + \bar{S}_{13} t_r + \bar{S}_{14} t_{r\theta} + \bar{S}_{15} t_{rz} + \bar{S}_{16} t_{\theta z}$$

$$v_{\theta,\phi} + \lambda^{\frac{1}{2}} v_r = \lambda^{\frac{1}{2}} (1 + \lambda y) \bar{S}_{21} t_z + \bar{S}_{22} t_\theta + \bar{S}_{23} t_r + \bar{S}_{24} t_{r\theta} + \bar{S}_{25} t_{rz} + \bar{S}_{26} t_{\theta z}$$

$$(1 + \lambda y) v_{\theta,x} + v_{z,\phi}$$

$$= \lambda^{\frac{1}{2}} (1 + \lambda y) \bar{S}_{61} t_z + \bar{S}_{62} t_\theta + \bar{S}_{63} t_r + \bar{S}_{64} t_{r\theta} + \bar{S}_{65} t_{rz} + \bar{S}_{66} t_{\theta z}$$

(4.2)

$$v_{r,y}^{(0)} = 0$$

$$v_{z,y}^{(1)} + v_{r,x}^{(0)} = 0$$

$$v_{\theta,y}^{(1)} + v_{r,\phi}^{(0)} = 0$$

$$v_{z,x}^{(1)} = S_{11}^{(0)} t_z^{(0)} + S_{12}^{(0)} t_{\theta}^{(0)} + S_{13}^{(0)} t_{\theta z}^{(0)}$$

$$v_{\theta,\phi}^{(1)} + v_r^{(0)} = S_{12}^{(0)} t_z^{(0)} + S_{22}^{(0)} t_{\theta}^{(0)} + S_{26}^{(0)} t_{\theta z}^{(0)}$$

$$v_{\theta,x}^{(1)} + v_{z,\theta}^{(1)} = S_{16}^{(0)} t_z^{(0)} + S_{26}^{(0)} t_{\theta}^{(0)} + S_{66}^{(0)} t_{\theta z}^{(0)}$$

$$t_{rz,y}^{(1)} + t_{z,x}^{(0)} + t_{\theta z,\phi}^{(0)} = 0$$

$$t_{r\theta,y}^{(1)} + t_{\theta,\phi}^{(0)} + t_{\theta z,x}^{(0)} = 0$$

$$t_{r,y}^{(2)} + t_{rz,x}^{(1)} + t_{r\theta,\phi}^{(1)} - t_{\theta}^{(0)} = 0$$

From the initial terms chosen for the first approximation system.

where v_r , v_θ , v_z are the components of displacement at $y = 0$ surface. You will find the linear y dependence of the in-plane displacements.

We will then obtain the in-plane stress-strain relations:

$$\begin{Bmatrix} t_z^{(0)} \\ t_\theta^{(0)} \\ t_{\theta z}^{(0)} \end{Bmatrix} = [C] \begin{Bmatrix} \epsilon_1 \\ \epsilon_2 \\ \epsilon_{12} \end{Bmatrix} + [C] \begin{Bmatrix} K_1 \\ K_2 \\ K_{12} \end{Bmatrix} y \quad (4.4)$$

where strain components ϵ_i and curvature components k_i defined at $y = 0$, surface for a conical shell are given as follows:

$$\begin{aligned} \epsilon_1 &= V_{z,x}^{(1)} & K_1 &= -V_{r,xx}^{(0)} / (\sin \alpha) \\ \epsilon_2 &= V_r^{(0)} (\sin \alpha) + V_{\theta,\phi}^{(1)} & K_2 &= -V_{r,\phi\phi}^{(0)} / (\sin \alpha) \\ \epsilon_{12} &= V_{\theta,x}^{(1)} + V_{z,\phi}^{(1)} & K_{12} &= -2V_{r,x\phi}^{(0)} / (\sin \alpha) \end{aligned}$$

(4.5)

and $[C]$ is the inverse of the first approximation compliance matrix.

$$[C] = \begin{bmatrix} s_{11}^{(0)} & s_{12}^{(0)} & s_{16}^{(0)} \\ s_{12}^{(0)} & s_{22}^{(0)} & s_{26}^{(0)} \\ s_{16}^{(0)} & s_{26}^{(0)} & s_{66}^{(0)} \end{bmatrix}^{-1} \quad (4.6)$$

The transverse stresses can now be solved by substituting the in-plane stresses obtained in (4.5) into the last equations of (4.7) and integrating with respect to y , we then get:

$$\begin{aligned} t_{rz}^{(1)} &= T_{rz}(x, \phi) - (A_{1j}\epsilon_{j,x} + B_{1j}K_{j,x}) - (A_{3j}\epsilon_{j,\phi} + B_{3j}K_{j,\phi}) \\ t_{r\theta}^{(1)} &= T_{r\theta}(x, \phi) - (A_{3j}\epsilon_{j,x} + B_{2j}K_{j,x}) - (A_{2j}\epsilon_{j,\phi} + B_{3j}K_{j,\phi}) \\ t_r^{(2)} &= T_r(x, \phi) - (T_{rz,x} + T_{r\theta,\phi})y + A_{2j}\epsilon_j + B_{2j}K_j + D_{1j}\epsilon_{j,xx} \\ &\quad + E_{1j}K_{j,xx} + 2D_{3j}\epsilon_{j,x} + 2E_{3j}K_{j,x} + D_{2j}\epsilon_{j,\phi} + E_{2j}K_{j,\phi} \end{aligned} \quad (4.7)$$

where T_{rz} , $T_{r\theta}$ and T_r are the stress components at $y = 0$ and A_{ij} , B_{ij} , D_{ij} and E_{ij} are the products obtained by integration of stress strain coefficient components over the thickness coordinate. The boundary conditions are:

$$\begin{aligned}
 \sigma_r = \tau_{rz} = \tau_{r\theta} = 0 \text{ at } r = a \text{ and} \\
 \sigma_r = p(\theta, z) \text{ when } r = a + h \\
 \text{also} \\
 \tau_{rz} = \tau_{r\theta} = 0 \text{ at } r = a + h
 \end{aligned}
 \tag{4.8}$$

Where $p^* = p / (\sigma\lambda)$, to non-dimensionalize.

Application

We assume here that the contact between layers is such that the strains are a continuous function in thickness coordinate. As the C_{ij} are piecewise continuous functions, the in-plane stresses are also continuous. We would expect the in-plane stresses to be discontinuous at the juncture of layers of dissimilar materials. The transverse stresses are continuous functions of the thickness coordinate. We assume here that the contact between layers is such that the strains are a continuous function in thickness coordinate. As the C_{ij} are piecewise continuous functions, the in-plane stresses are also continuous. We would expect the in-plane stresses to be discontinuous at the juncture of layers of dissimilar materials. The transverse stresses are continuous functions of the thickness coordinate.

Although, as mentioned above, the theory developed can consider. Random layering, numerical results are to be carried out for a four-layer symmetric angle ply configuration. For this configuration the angle of elastic axes γ is oriented at $+\gamma, -\gamma, -\gamma, +\gamma$ with the shell axis and the layers are of equal thickness.

The cone is taken to be clamped at both ends but free to rotate and extend axially at one end. We now choose it to be such that there exists no coupling between N_z and K_z and M_z and ϵ_{id} .

As the loading applied is axisymmetric, all the stresses and strains are also taken to be axisymmetric.

Numerical calculations are now carried out for a shell of the following dimensions:

We will then obtain $\lambda = 0.025$ ($0.1 / 4.0$). Each of the layers is taken to be 0.025 inches thick and thus the dimensionless distances from the bottom of the first layer are given by

$$S_1 = 0 \quad S_2 = 0.25 \quad S_3 = 0.5 \quad S_4 = 0.75 \quad S_5 = 1.0$$

Each layer of the symmetric angle ply configuration (elastic symmetry axes y are oriented at $(+y, -y, -y, +y)$) is taken to be orthotropic with engineering elastic coefficients of Apollo 38-750 Graphite Fiber as follows:

$$E_1 = 300 \text{ Gpa}$$

$$E_2 = 260 \text{ Gpa}$$

$$G_{12} = 52 \text{ Gpa}$$

Here, direction 1 signifies the direction parallel to the fibers, while 2 is the transverse direction. Angles chosen were $\gamma = 0, 15, 30, 45$ and 60 . Use of the transformation equations (2.7) then yields the mechanical properties for the different symmetric angle ply configurations.

Shown in Fig. 3 is the variation of the dimensionless radial displacement with the actual distance along the axis for the different ply angles. The reference surface for the chosen configuration is given by $d/h = 1/2$. The integration constants are determined from the edge conditions. The radial deformations of the length scales $(Lh)^{1/2}$ is shown in Fig. 1 and 2. One will notice the radial deformation rapidly increases along the length scale near the edge while further inside deformations show little change. This area is the so-called edge effect zone or boundary layer; the American Petroleum institute adopted similar theories for its cylindrical shell gas storage tank design (API 650). Locating the edge effect zone can simplify the calculation of circumferential and longitudinal stiffeners as shown in the References [8] through [11], in relation to the buckling problems, Reference [15].

It is also shown that wide variations in the magnitude of radial displacement take place with change in the cross-ply angle. The maximum displacement occurs at $\gamma = 30$ degrees while the minimum displacement is at $\gamma = 60$ degrees.

In each case, the displacements increase in γ up to $\gamma = 30$ degrees, and thereafter will decrease. Weak and smooth edge effects are the case for large cross-ply angles.

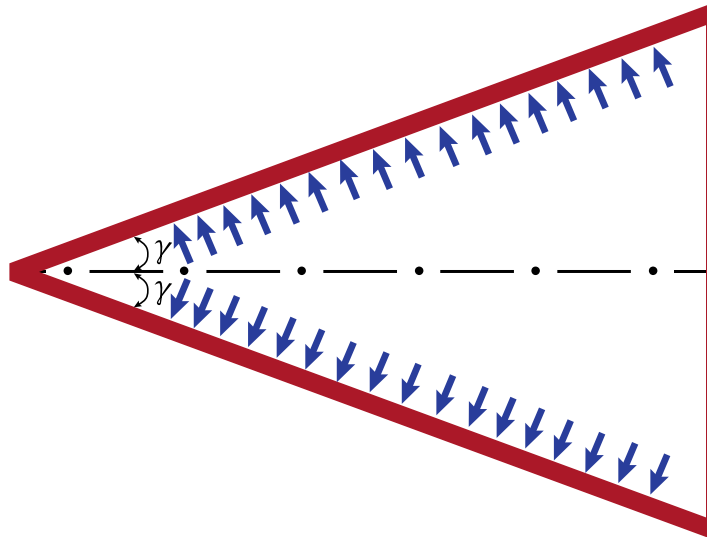


Figure (3) Conical Shell under Internal Pressure

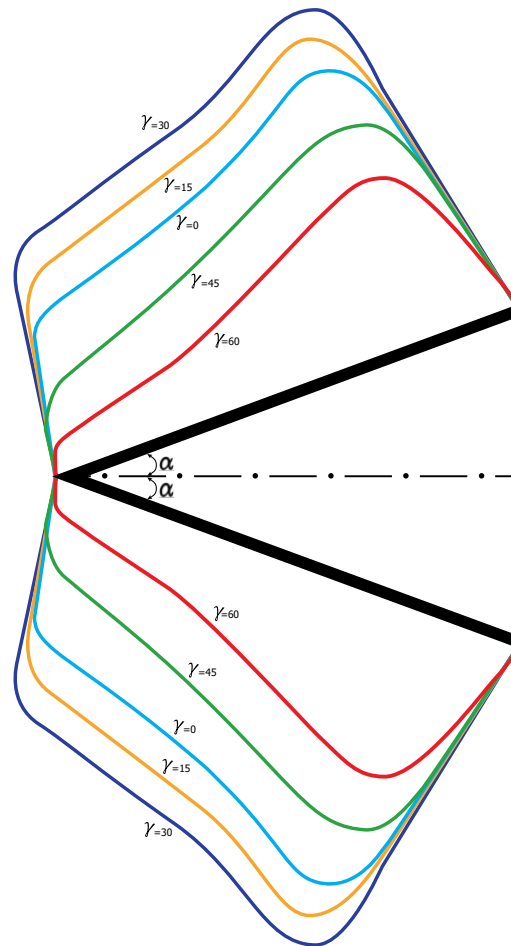


Figure (4) Axi-Symmetric Deformations

CONCLUSIONS

In this paper, we first discuss how the approximation shell theories are derived by use of the method of asymptotic integration of the exact three-dimensional elasticity equations for a hybrid anisotropic conical shell.

The analysis remains valid for materials, which are non-homogeneous to the extent that their properties are allowed to vary with the thickness coordinate.

As a result of the application of this method, one can obtain shell approximate theory of various orders in a systematic manner. The first approximation theories derived in this paper represent the simplest possible shell theories for the corresponding length scales considered. Although 21 elastic coefficients are present in the original formulation of the problem, only six appear in the first approximation theories. The shell theories thus assume the existence of a plane of elastic symmetry.

It was seen that various shell theories are obtained by using different combinations of the length scales introduced in the non-dimensionalization of the coordinates and that each theory possesses unique properties such as the orders of stress magnitudes, displacement components and edge effect penetration. In this research, we considered a theory with length scales of $(Lh)^{1/2}$ for both longitudinal and circumferential directions, where L the actual cone length and h is the total thickness of the cone.

To illustrate the use of these theories, the application to layered shells was shown. The results of the analysis showed the radial displacement first increases with increase of the angle between the axis of elastic symmetry and longitudinal shell axis, being largest at 30 degrees. The theory of axial and circumferential length scale of $(Lh)^{1/2}$ show as a significant edge effect exists and that the penetration of the edge effect changes with the angle in similar fashion as for the radial displacement, being deepest at 30 degrees.

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REFERENCES

- [1] Love, A. E. H., A treatise on the mathematical theory of elasticity, 4th edition. New York, NY, USA: Dover Publications, 2011.
- [2] Donnell, L. H., Beams, Plates and Shells (Engineering societies monographs): McGraw-Hill Inc.: ISBN-13: 9780070175938.
- [3] Reissner, E., "The effect of transverse shear deformation on the bending of elastic plates," J. Appl. Mech., vol 12, no A, pp. 69-77, 1945.
- [4] Johnson, M. W. and Widera, O. E., "An asymptotic dynamic theory for cylindrical shells," Studies Appl. Math., vol 48, pp. 205, 1969.
- [5] Widera, O. E., "An asymptotic theory for the motion of elastic plates," Acta mechanica, vol

- 9, no 54, 1970.
- [6] Vlasov, V. Z., "General theory of shells and its application in engineering", NASA TT F-99, National Tech. Information Service.
- [7] De Renzo, D.J., "Advanced Composite Materials, Products and Manufactures", Noyes Data Corporation, ISBN 0-8155-1155-8.
- [8] Heslehurst, R. B., "Design and analysis of structural joints with composite materials", DES. "Details of Semi-Membrane Shell Theory of Hybrid Anisotropic Materials" International Journal of Composite Materials 2018.8(3):47-56, Chung, S.W., Hong, S.G. and Ju, G.S.,
- [9] "Pseudo –Membrane Shell Theory of Hybrid Anisotropic Materials", Journal of Composite Structures, Volume 160, 15 January 2017, Chung, S.W., Hong, S.G.
- [10] "Details of Semi-Membrane Shell Theory of Hybrid Anisotropic Materials" International Journal of Composite Materials 2018.8(3):47-56, Chung, S.W., Hong, S.G. and Ju, G.S.,
- [11] "A Comparison of Membrane Shell Theories of Hybrid Anisotropic Materials", European Journal of Engineering and Technology, Volume 4, No.5,2016
<http://www.idpublications.org/ejet-vol-4-no-5-2016/>
- [14] "Pure Membrane, Pseudo Membrane, and Semi Membrane Shell Theories of Hybrid Anisotropic Materials", Journal of Material Science and Engineering A 8 (5-6) (2018) 121-135, Chung, S.W., Hong, S.G., Ju, G.S.
- [15] Semi-Membrane and Effective Length Theory of hybrid Anisotropic Materials, Samuel Chung, Gi Su Ju. Submission Date: 05/23/2017(MM/DD/YYYY) International Journal of Composite Materials. volume 7 (3), 2017: Contact Us: editor@sapub.org , ID:110900221
- [16] T. Tokugawa, "Experiments on the elastic stability of a thin-wall cone under uniform normal pressure on all sides, and an approximate method for computing its collapsible pressure", in Congress of Applied Mechanics League, Zōsen Kyōkai (Shipbuilding Association), Miscellaneous Publications, 125, 1932, 151-169.
- [17] C. T. F. Ross, "Pressure Vessels: External Pressure Technology", Horwood Publishing Ltd, Chichester UK, 2001.
- [18] P. Seide, "Axisymmetrical buckling of circular cones under axial compression", Journal of Applied Mechanics, Transaction of ASME, 23(4), 1956, 625-628.
- [19] P. Seide, "Buckling of circular cones under axial compression", Journal of Applied Mechanics, 28(2), 1961, 315-326. [5] J. Singer, "Buckling of circular conical shells under uniform axial compression", AIAA Journal, 3 (5),
- [20] S. Kobayashi, "The influence of prebuckling deformation on the buckling load of truncated conical shells under axial compression", NASA CR-707, 1967, pp. 1-41
- [21] J. Tani, "Influence of prebuckling deformations on the buckling of truncated conical shells under axial compression", Transactions of Japan Society for Aeronautical and Space Sciences, 16(34), 1973, 232-245. [8] J. Tani, N. Yamaki, "Buckling of truncated conical shells under axial compression", AIAA Journal, (3), 1970, 568-571. [9] M. Baruch, O. Harari, J. Singer, "Low buckling loads of axially compressed conical shells", Journal of Applied Mechanics, Transactions of the ASME, 37(2), 1970, 384-392.
- [23] C. H. Chang, L. Katz, "Buckling of axially compressed conical shells", Journal of Engineering, Mechanics Division, 106(3), 1980, 501–516.